

# Holographic transform : the Rankin-Cohen case

Labriet Quentin

Université de Reims

Sochi, February 25, 2020

# Representation theory

## Branching laws

Let  $G$  be a group and  $G'$  a subgroup of  $G$ . We consider an irreducible unitary representation  $\pi$  of  $G$  and we are interested in the problem of the restriction of  $\pi$  to  $G'$ , and the decomposition of  $\pi$  into irreducible representations  $\rho_\lambda$  of  $G'$ .

# Representation theory

## Branching laws

Let  $G$  be a group and  $G'$  a subgroup of  $G$ . We consider an irreducible unitary representation  $\pi$  of  $G$  and we are interested in the problem of the restriction of  $\pi$  to  $G'$ , and the decomposition of  $\pi$  into irreducible representations  $\rho_\lambda$  of  $G'$ .

Branching law : multiplicity free case

$$\pi|_{G'} \simeq \sum_{\lambda \in \hat{G}'} m_\lambda \rho_\lambda \text{ with } m_\lambda \in \{0; 1\}$$

# Representation theory

## Branching laws

Let  $G$  be a group and  $G'$  a subgroup of  $G$ . We consider an irreducible unitary representation  $\pi$  of  $G$  and we are interested in the problem of the restriction of  $\pi$  to  $G'$ , and the decomposition of  $\pi$  into irreducible representations  $\rho_\lambda$  of  $G'$ .

### Branching law : multiplicity free case

$$\pi|_{G'} \simeq \sum_{\lambda \in \hat{G}'} m_\lambda \rho_\lambda \text{ with } m_\lambda \in \{0; 1\}$$

Multiplicity free branching law

$\Leftrightarrow$

$$\dim \operatorname{Hom}_{G'}(\pi|_{G'}, \rho_\lambda) = 1$$

# Representation theory

## Symmetry breaking and holographic transform

In this context, a non zero element  $T_\lambda \in \text{Hom}_{G'} (\pi|_{G'}, \rho_\lambda)$  is called a *symmetry breaking operator*, and the collection  $(T_\lambda)_{\lambda \in \hat{G}'}$  is called a *symmetry breaking transform*.

# Representation theory

## Symmetry breaking and holographic transform

In this context, a non zero element  $T_\lambda \in \text{Hom}_{G'}(\pi|_{G'}, \rho_\lambda)$  is called a *symmetry breaking operator*, and the collection  $(T_\lambda)_{\lambda \in \hat{G}'}$  is called a *symmetry breaking transform*.

One can reverse the arrows in the previous definition, and a non zero element  $P_\lambda \in \text{Hom}_{G'}(\rho_\lambda, \pi|_{G'})$  is called an *holographic operator*, and the collection  $(P_\lambda)_{\lambda \in \hat{G}'}$  is called an *holographic transform*.

# Representation theory

## Symmetry breaking and holographic transform

In this context, a non zero element  $T_\lambda \in \text{Hom}_{G'}(\pi|_{G'}, \rho_\lambda)$  is called a *symmetry breaking operator*, and the collection  $(T_\lambda)_{\lambda \in \hat{G}'}$  is called a *symmetry breaking transform*.

→ Spectral analysis.

One can reverse the arrows in the previous definition, and a non zero element  $P_\lambda \in \text{Hom}_{G'}(\rho_\lambda, \pi|_{G'})$  is called an *holographic operator*, and the collection  $(P_\lambda)_{\lambda \in \hat{G}'}$  is called an *holographic transform*.

→ Spectral synthesis.

# Geometric setting

Holomorphic discrete series representation for  $SL_2(\mathbb{R})$

Let  $\lambda \in \mathbb{N} \setminus \{0; 1\}$ , and define the weighted Bergmann space on the Poincaré upper half-plane :

$$H_\lambda^2(\Pi) = \mathcal{O}(\Pi) \cap L^2(\Pi, y^{\lambda-1} dx dy)$$

We then define the holomorphic discrete series representation  $\pi_\lambda$  of  $SL_2(\mathbb{R})$  for  $f \in H_\lambda^2(\Pi)$  by the formula

$$\pi_\lambda(g)f(z) = (cz + d)^{-\lambda} f\left(\frac{az + b}{cz + d}\right)$$

for  $g^{-1} = \begin{pmatrix} a & b \\ c & d \end{pmatrix}$ .



We consider the case  $(G, G') = (SL_2(\mathbb{R}) \times SL_2(\mathbb{R}), \Delta SL_2(\mathbb{R}))$ .

We are interested in the tensor product representation of two discrete series  $\pi_{\lambda'} \otimes \pi_{\lambda''}$  which is considered as a representation of  $SL_2(\mathbb{R})$  on the space  $H_{\lambda'}^2(\Pi) \hat{\otimes} H_{\lambda''}^2(\Pi)$ .

## Branching law

$$\pi_{\lambda'} \otimes \pi_{\lambda''}|_{SL_2(\mathbb{R})} \simeq \sum_{l \in \mathbb{N}} \pi_{\lambda' + \lambda'' + 2l}$$

Multiplicity free branching law

$$\Leftrightarrow$$

$$\dim \operatorname{Hom}_{SL_2(\mathbb{R})}(\pi_{\lambda'} \otimes \pi_{\lambda''}|_{SL_2(\mathbb{R})}, \pi_{\lambda'''}) = 1$$

# Geometric setting

## Rankin-Cohen transform

Let  $\lambda', \lambda'', \lambda''' \in \mathbb{N} \setminus \{0; 1\}$  such that  $l = \frac{1}{2}(\lambda''' - \lambda' - \lambda'') \in \mathbb{N}$ .

For  $f \in H_{\lambda'}^2(\Pi) \hat{\otimes} H_{\lambda''}^2(\Pi)$ , we define the Rankin-Cohen brackets as :

### Rankin-Cohen bi-differential operators

$$RC_{\lambda', \lambda''}^{\lambda'''}(f)(z) = \sum_{j=0}^l \frac{(-1)^j (\lambda' + l - j)_j (\lambda'' + j)_{l-j}}{j!(l-j)!} \frac{\partial^l f}{\partial z_1^{l-j} \partial z_2^j}(z, z)$$

Each operator  $RC_{\lambda', \lambda''}^{\lambda'''}$  is a *symmetry breaking operator* from  $H_{\lambda'}^2(\Pi) \hat{\otimes} H_{\lambda''}^2(\Pi)$  to  $H_{\lambda'''}^2(\Pi)$ , and the family of all  $RC_{\lambda', \lambda''}^{\lambda'''}$  defines a *symmetry breaking transform*.

# Holographic inversion

## Relative reproducing kernel

We define  $K_{\lambda', \lambda''}^{\lambda'''}$ , called the relative reproducing kernel, for  $w_2, w_1, z \in \Pi$  :

$$K_{\lambda', \lambda''}^{\lambda'''}(z, w_1, w_2) = (w_2 - w_1)^l \left( \frac{w_1 - \bar{z}}{2i} \right)^{-(\lambda' + l)} \left( \frac{w_2 - \bar{z}}{2i} \right)^{-(\lambda'' + l)}$$

## Theorem (L.)

Set  $\lambda', \lambda'', \lambda''' \in \mathbb{N} \setminus \{0; 1\}$  such that  $l = \frac{1}{2}(\lambda''' - \lambda' - \lambda'') \in \mathbb{N}$ . Let  $w_1, w_2 \in \Pi$ , and  $g \in H_{\lambda'''}^2(\Pi)$ . Then the adjoint operator  $(RC_{\lambda', \lambda''}^{\lambda'''})^*$  is given by :

$$(RC_{\lambda', \lambda''}^{\lambda'''})^* g(w_1, w_2) = C \int_{\Pi} g(z) K_{\lambda', \lambda''}^{\lambda'''}(z, w_1, w_2) d\mu(z)$$

And this is an holographic operator.

# Holographic inversion

An integral operator

Set  $\lambda', \lambda'', \lambda''' \in \mathbb{N} \setminus \{0; 1\}$  such that  $l = \frac{1}{2}(\lambda''' - \lambda' - \lambda'') \in \mathbb{N}$ . Define the operator  $\Psi_{\lambda', \lambda''}^{\lambda'''}$  on the space  $H_{\lambda'''}(\Pi)$  by

$$\Psi_{\lambda', \lambda''}^{\lambda'''}(g)(w_1, w_2) = \frac{(w_1 - w_2)^l}{2^{\lambda' + \lambda'' + 2l - 1} l!} \int_{-1}^1 g(w(v)) (1-v)^{\lambda' + l - 1} (1+v)^{\lambda'' + l - 1} dv$$

## Theorem (Kobayashi, Pevzner)

*The operator  $\Psi_{\lambda', \lambda''}^{\lambda'''}$  is an holographic operator from  $H_{\lambda'''}(\Pi)$  to  $H_{\lambda'}^2(\Pi) \hat{\otimes} H_{\lambda''}^2(\Pi)$ .*

For  $\lambda \in \mathbb{N} \setminus \{0.1\}$ , we consider the space  $L^2(\mathbb{R}^+, t^{1-\lambda} dt) := L^2_\lambda(\mathbb{R}^+)$ . Then the Fourier transform defined, for  $f \in L^2_\lambda(\mathbb{R}^+)$ , by :

$$\mathcal{F}f(z) = \int_{\mathbb{R}^+} f(t) e^{itz} dt$$

is an one-to-one isometry (up to a scalar) from  $L^2_\lambda(\mathbb{R}^+)$  to  $H^2_\lambda(\Pi)$ . This allows us to transport the holomorphic discrete series representation of  $SL_2(\mathbb{R})$  on this space.

We define the counterpart of  $RC_{\lambda',\lambda''}^{\lambda'''}$  in this model thanks to the following commutative diagram

$$\begin{array}{ccc}
 L_{\lambda'}^2(\mathbb{R}^+) \hat{\otimes} L_{\lambda''}^2(\mathbb{R}^+) & \xrightarrow{\widehat{RC_{\lambda',\lambda''}^{\lambda'''}}} & L_{\lambda'''}^2(\mathbb{R}^+) \\
 \mathcal{F} \otimes \mathcal{F} \downarrow & & \downarrow \mathcal{F} \\
 H_{\lambda'}^2(\Pi) \hat{\otimes} H_{\lambda''}^2(\Pi) & \xrightarrow{RC_{\lambda',\lambda''}^{\lambda'''}} & H_{\lambda'''}^2(\Pi)
 \end{array}$$

Symmetry breaking operator in the holomorphic and  $L^2$ -model

### Theorem (Kobayashi, Pevzner)

Set  $\lambda', \lambda'', \lambda''' \in \mathbb{N} \setminus \{0; 1\}$  such that  $l = \frac{1}{2}(\lambda''' - \lambda' - \lambda'') \in \mathbb{N}$ . Let  $F \in L^2_{\lambda', \lambda''}(\mathbb{R}^+ \times \mathbb{R}^+)$  then

$$\widehat{RC^{\lambda'''}_{\lambda', \lambda''}} F(t) = \frac{t^{l+1}}{2i^l} \int_{-1}^1 P_l^{\lambda'-1, \lambda''-1}(v) F\left(\frac{t}{2}(1-v), \frac{t}{2}(1+v)\right) dv$$

This operator is a *symmetry breaking operator* in the  $L^2$ -model.

For the holographic transform we have the following situation

$$\begin{array}{ccc} L^2_{\lambda'}(\mathbb{R}^+) \hat{\otimes} L^2_{\lambda''}(\mathbb{R}^+) & \longleftarrow & L^2_{\lambda'''}(\mathbb{R}^+) \\ \mathcal{F} \otimes \mathcal{F} \downarrow & & \downarrow \mathcal{F} \\ H^2_{\lambda'}(\Pi) \hat{\otimes} H^2_{\lambda''}(\Pi) & \xleftarrow{\Psi_{\lambda', \lambda''}^{\lambda'''}} & H^2_{\lambda'''}(\Pi) \end{array}$$

Holographic operator in the holomorphic and  $L^2$ -model



### Theorem (Kobayashi, Pevzner)

Set  $\lambda', \lambda'', \lambda''' \in \mathbb{N} \setminus \{0; 1\}$  such that  $l = \frac{1}{2}(\lambda''' - \lambda' - \lambda'') \in \mathbb{N}$ . We define the following operator which associate to a function  $g(t)$  defined on  $\mathbb{R}^+$  a function of two variables on  $\mathbb{R}^+ \times \mathbb{R}^+$  :

$$\Phi_{\lambda', \lambda''}^{\lambda'''} g(x, y) = \frac{x^{\lambda'-1} y^{\lambda''-1}}{(x+y)^{\lambda'+\lambda''+l-1}} P_l^{\lambda'-1, \lambda''-1} \left( \frac{y-x}{x+y} \right) g(x+y)$$

Then  $\Phi_{\lambda', \lambda''}^{\lambda'''}$  is an holographic operator from  $L_{\lambda'''}^2(\mathbb{R}^+)$  to  $L_{\lambda', \lambda''}^2(\mathbb{R}^+ \times \mathbb{R}^+)$

# A geometric interpretation

A stratification of the cone  $\mathbb{R}^+ \times \mathbb{R}^+$

We use the following diffeomorphism  $\theta$  from  $\mathbb{R}^+ \times (-1, 1)$  to  $\mathbb{R}^+ \times \mathbb{R}^+$  :

$$\theta(t, v) = \left( \frac{t}{2}(1 - v), \frac{t}{2}(1 + v) \right)$$

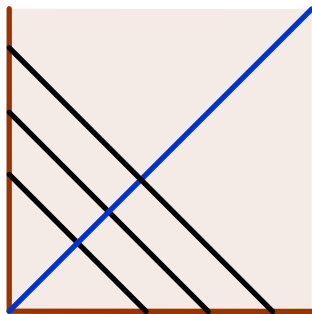
# A geometric interpretation

A stratification of the cone  $\mathbb{R}^+ \times \mathbb{R}^+$

We use the following diffeomorphism  $\theta$  from  $\mathbb{R}^+ \times (-1, 1)$  to  $\mathbb{R}^+ \times \mathbb{R}^+$  :

$$\theta(t, v) = \left( \frac{t}{2}(1 - v), \frac{t}{2}(1 + v) \right)$$

This corresponds to the following stratification of the cone  $\mathbb{R}^+ \times \mathbb{R}^+$



# A geometric interpretation

Using this diffeomorphism, one can show the following isomorphism

$$\begin{aligned} & L^2_{\lambda', \lambda''}(\mathbb{R}^+ \times \mathbb{R}^+) \\ \simeq & L^2\left(\mathbb{R}^+, t^{\lambda' + \lambda'' - 1} dt\right) \hat{\otimes} L^2\left((-1; 1), (1 - v)^{\lambda' - 1} (1 + v)^{\lambda'' - 1} dv\right) \end{aligned}$$

# A geometric interpretation

Using this diffeomorphism, one can show the following isomorphism

$$\begin{aligned} & L^2_{\lambda', \lambda''}(\mathbb{R}^+ \times \mathbb{R}^+) \\ \simeq & L^2\left(\mathbb{R}^+, t^{\lambda' + \lambda'' - 1} dt\right) \hat{\otimes} L^2\left((-1; 1), (1 - v)^{\lambda' - 1} (1 + v)^{\lambda'' - 1} dv\right) \\ \simeq & \bigoplus_{l \geq 0} L^2\left(\mathbb{R}^+, t^{\lambda' + \lambda'' - 1} dt\right) \hat{\otimes} \mathbb{C} \cdot P_l^{\lambda' - 1, \lambda'' - 1}(v) \end{aligned}$$

# A geometric interpretation

Using this diffeomorphism, one can show the following isomorphism

$$\begin{aligned} & L^2_{\lambda', \lambda''}(\mathbb{R}^+ \times \mathbb{R}^+) \\ & \simeq L^2\left(\mathbb{R}^+, t^{\lambda' + \lambda'' - 1} dt\right) \hat{\otimes} L^2\left((-1; 1), (1 - v)^{\lambda' - 1} (1 + v)^{\lambda'' - 1} dv\right) \\ & \simeq \bigoplus_{l \geq 0} L^2\left(\mathbb{R}^+, t^{\lambda' + \lambda'' - 1} dt\right) \hat{\otimes} \mathbb{C} \cdot P_l^{\lambda' - 1, \lambda'' - 1}(v) \end{aligned}$$

## Theorem (L.)

*The orthogonal projection on the space  $L^2\left(\mathbb{R}^+, t^{\lambda' + \lambda'' - 1} dt\right) \hat{\otimes} \mathbb{C} \cdot P_l^{\lambda' - 1, \lambda'' - 1}(v)$  is a symmetry breaking operator for the  $SL_2(\mathbb{R})$  representation transported from the  $L^2$ -model.*

Thank you for your attention.